RMS Delay and Coherence Bandwidth Measurements in Indoor Radio Channels in the UHF Band

Mercedes Sánchez Varela and Manuel García Sánchez, Member, IEEE

Abstract—A study of time dispersion in different indoor line-of-sight radio channels in the 492–862 MHz band is presented in this paper. A combined method to filter the noise in the measured impulse response is described. The effect of frequency windowing on the impulse responses and the root mean square (rms) delay spread is also investigated. It has been found that, in general, the use of windows with lower side-lobe levels yields larger values of the rms delay spread. The relation between the mean delay and the rms delay spread has also been studied for copolar and crosspolar channels. The dependence of the coherence bandwidth on the rms delay spread has been considered, and an inverse relation has been tested for both components.

Index Terms—Communication channels, communication systems, parameter estimation, UHF measurements, UHF radio propagation.

I. INTRODUCTION

The dramatic increase of the number of digital radio communication systems within buildings has demanded wide band characterization of indoor radio channels. Signal time dispersion is one of the main study issues because it limits the maximum symbol rate that can be used without intersymbol interference [1]. The impulse response (IR) of the radio channel and other parameters, as the mean delay ($\tau_{\text{mean}}$) and the root mean square (rms) delay spread ($\tau_{\text{rms}}$), are frequently used to characterize time dispersion of the channel. The coherence bandwidth ($B_c$), a parameter closely related with $\tau_{\text{rms}}$, is also used to describe frequency selectivity in the radio channel.

Previous works have investigated the relation between $B_c$ and $\tau_{\text{rms}}$, and different solutions have been found for outdoor environments [1]–[6]. But at present, few experimental results exist that describe the relation between $B_c$ and $\tau_{\text{rms}}$ in indoor environments. In this paper, an experimental study on the relation between these parameters inside buildings is presented. The results are based on wide-band measurements taken in three different indoor environments with a frequency swept radio channel sounder. The sounder, based on a vector network analyzer (VNA), has been used to measure the frequency response of the radio channel in the 470–862 MHz frequency band [7]. The measurement setup is shown in Fig. 1. The copolar component was measured with two vertical 3-dBi omnidirectional dipole antennas connected to the input and output ports of the VNA and placed on the top of a 1-m-height tripod. To measure the crosspolar component, the vertical dipole at the receiver was replaced by a 5-dBi omnidirectional antenna consisting of two orthogonal folded dipoles. This antenna was horizontally polarized.

The measurement system output is the frequency response of the devices connected between its two ports, including the channel, antennas, cables, amplifiers, and frequency response of the VNA itself. To compensate the effect of the system on the measurements, a calibration was carried out before the radio channel measurements. With this purpose, the antennas were removed so that the transmitter output was directly connected to the receiver input, its frequency response being measured. This result was automatically subtracted from the subsequent channel measurements, thus reducing the effect of the system on the measurement

$$H(f, t)_{\text{channel}} = \frac{H(f, t)_{\text{measured}}}{H(f)_{\text{system}}}. \quad (1)$$

If the channel is linear, the IR can be calculated from the frequency response using the inverse Fourier transform [8]

$$h(\tau, t) = \int_{-\infty}^{\infty} H(f, t)e^{j2\pi ft} df. \quad (2)$$

However, since the measurements were performed over a limited frequency band, there was a windowing effect on the results

$$h(\tau, t)_{\text{estimated}} = \int_{f_{\text{min}}}^{f_{\text{max}}} H(f, t)e^{j2\pi ft} df = \int_{-\infty}^{\infty} H_{\text{channel}}(f, t) \cdot H_{\text{window}}(f)e^{j2\pi ft} df. \quad (3)$$
The shape of the window used in the frequency measurements affects the IRs as follows:

\[ h(\tau,t)_{\text{estimated}} = h(\tau,t)_{\text{channel}} \ast h_{\text{window}}(\tau). \]  

As a result, the impulse response is averaged and the delay resolution is reduced. Four different windows were considered in this study: rectangular, Hanning, Blackman–Harris, and Kaiser–Bessel. The side-lobe levels (SLLs) for these windows are given in Table I. Lower SLLs yield wider main lobes [9]. A wider main lobe results in lower delay resolution in the IR.

If the channel satisfies the wide sense stationary uncorrelated scattering (WSSUS) assumption, the power delay profile (PDP) is given by

\[ P_{h}(\tau) = \langle \lvert h(\tau,t) \rvert^{2} \rangle. \]  

Wide-band parameters, such as the mean delay \( \tau_{\text{mean}} \), the rms delay spread \( \tau_{\text{rms}} \), and the coherence bandwidth at level \( \rho \in [0,1] \), \( B_{\rho} \) are calculated. The mean delay \( \tau_{\text{mean}} \) is the average of the delays of all paths

\[ \tau_{\text{mean}} = \frac{\int_{0}^{\tau_{\text{max}}} \tau P_{h}(\tau) d\tau}{\int_{0}^{\tau_{\text{max}}} P_{h}(\tau) d\tau}. \]  

The radio channel time dispersion is characterized by the rms delay spread \( \tau_{\text{rms}} \), calculated as

\[ \tau_{\text{rms}} = \sqrt{\frac{\int_{0}^{\tau_{\text{max}}} (\tau - \tau_{\text{mean}})^{2} P_{h}(\tau) d\tau}{\int_{0}^{\tau_{\text{max}}} P_{h}(\tau) d\tau}}. \]  

From (4), it can be demonstrated [10] that the channel delay spread is overestimated, the increase being due to the window shape

\[ (\tau_{\text{rms}})_{\text{estimated}}^{2} = (\tau_{\text{rms}})_{\text{channel}}^{2} + (\tau_{\text{rms}})_{\text{window}}^{2} \]  

where \( (\tau_{\text{rms}})_{\text{window}} \) is the delay spread corresponding to the window impulse response. Values of \( (\tau_{\text{rms}})_{\text{window}} \) for the four different windows considered in this work have been calculated and are presented in Table I.

The frequency correlation function of the radio channel can be obtained from the PDP as

\[ R_{T}(\Delta f) = \int_{-\infty}^{\infty} P_{h}(\tau) e^{-j2\pi \Delta f \tau} d\tau. \]

For a particular correlation level \( \rho \), typically 0.9, 0.7, or 0.5, \( B_{\rho} \) is the minimum frequency separation for which the norm of the frequency correlation function crosses this level. As an example, \( B_{0.5} \) is calculated as

\[ B_{0.5} = \min(\Delta f) \text{ such that } |R_{T}(\Delta f)| = 0.5. \]

For low values of \( \rho \), this parameter represents the minimum frequency separation to have the components of the radio signal sufficiently uncorrelated.

### III. Environment Description

Measurements were carried out in three different environments: a computer laboratory with rows of PC desks; an electronics laboratory with electronic equipments such as oscilloscopes, signal generators, and synthesizers on benches; and a corridor. The plan of these environments, with the dimensions in centimeters, is given in Fig. 2. In each environment, the receiver antenna was shifted to various positions separated one-eighth of a wavelength along a line. A total of 260, 240, and 50 different positions were used in the corridor, the computer laboratory, and the electronics laboratory, respectively. The transmitter and receiver antenna positions are also shown in Fig. 2. As can be seen, the LoS condition is always satisfied.

### IV. Impact of Windowing on the Estimated Power Delay Profiles and Delay Spreads

Despite the fact that the measurement system averages the results of several frequency sweeps to reduce the noise in the measurement, the frequency response and the IR will be corrupted by noise. Generally, a noise threshold is applied to raw IRs in order to separate actual multipath components from noise. Several methods to calculate this threshold noise can be found in
In the literature [11]–[13], an empirical/analytical method is proposed that is not affected by the actual noise level in the environment. To find the dynamic range noise floor, first the lowest 25% of the delay profile amplitude points are sorted, and then the highest and the lowest 25% of these low amplitude points are removed (median filtering). The remaining low amplitude points are averaged to yield a power level that is the dynamic range noise floor. In this method, we assume an average signal-to-noise ratio of at least 10 dB across the measured profile. In [12], it is assumed that the noise level is much lower than the side-lobe level (SLL) of the rectangular window, so the noise threshold is given by the SSL. Other authors, such as [13], take into account the thermal noise of the system so that the threshold becomes a function of the noise level. In this work, a new combined method to evaluate the noise threshold has been developed and used. Three different levels are calculated, and the threshold is chosen as the maximum of these three levels.

The calculation of the first level is based on measurements of the actual noise. Real and imaginary parts of the noise are assumed to be random variables that follow a zero-mean Gaussian distribution. From the corresponding measured cumulative distribution functions (cdfs) of real and imaginary parts of the noise, the value below which the random variables remain the 99.999% of the time is determined. The level is taken as the square root of the sum of both squared values. The second level is also calculated from measurements of the noise, but it becomes from the noise amplitude, which is assumed to follow a Rayleigh distribution [14]. The measured cdf of noise amplitude is calculated and the limit for 99.999% of the time is selected as the level. The third level is calculated as SSL (dB) below the peak signal level. This level does not depend on the
actual noise level but on the frequency window used for the measurements. The noise threshold estimation is improved by the proposed method because it combines the limits given by the actual noise power and the limit imposed by the window SLL.

To illustrate the effect of the frequency windowing on the IRs, the PDPs and the delay parameters have been calculated, with different frequency windows, at a fixed position in the corridor. They are shown in Fig. 3. The \( \tau_{\text{MEAN}} \) values range from 5.5 to 24.3\,ns, depending on the window used. In view of the values of \( (\tau_{\text{MEAN}})^{\text{window}} \) of Table I, the large variation of the delay spread cannot be completely explained by the only effect described in (8).

In order to gain a better understanding of the window effect on the delay spread, the threshold used to distinguish signal from noise in the PDPs has to be studied. When the rectangular window is used (SSL = -13\,dB), the overall threshold is fixed by the window SLL, because the two other levels are much lower. In this case, due to the high value of the threshold, some low-power multipath components are not considered and \( \tau_{\text{MEAN}} \) is underestimated. The other windows present lower SLLs, so the threshold is given by the actual noise level. For these cases, a larger number of multipath components are taken into account in the calculation and a better estimation of the delay spread is obtained. In general, lower window SLLs yield larger values of \( \tau_{\text{MEAN}} \). In terms of signal delay dispersion, this means that a larger value is obtained when windows such as Blackman–Harris are used.

V. RMS DELAY AND MEAN DELAY

The dependence of the rms delay on the mean delay has been analyzed in the three different environments described in Section III for copolar (vertical–vertical) and crosspolar (vertical–horizontal) components.

From the measurements collected in the three environments, \( \tau_{\text{MEAN}} \) and \( \tau_{\text{RMS}} \) have been estimated. They are plotted in Figs. 4, 7, and 9 as a function of the separation between antennas. The values obtained are comparable to those found in the literature for indoor environments [10], [12], [15], [16]. The mean delay
exhibits a trend to increase with distance. Some authors have reported that rms delay spread increases with antenna separation [15]. However, the rms delay spread is almost constant in the plots presented in this paper because of the short receiver antenna path. As the transmitter antenna is moved away from the receiver antenna, the amplitudes of the reflected signals relative to the direct path become larger, and this produces the increase of the mean and rms delays. However, due to the limited delay resolution of the measurement system, some oscillations are superimposed to the trend described above. After a thorough analysis of the IRs, it has been found that contributions with close delays cannot be resolved and interfere, giving rise both to a fast spatial variation of some IR and PDP components and to the oscillation of the delay parameters with the distance between antennas. It should be noted that no averaging of several squared envelopes of the impulse responses is performed because the effect of fast variation is also an issue of study. This explains the fluctuations that the estimated mean delays and rms delay spreads exhibit as the separation between the transmitter and the receiver is changed.

Values of both parameters are lower for the copolar than for the crosspolar component because the copolar IRs present a dominant ray while the crosspolar ones do not. To illustrate these differences, a sample of the IRs corresponding to both components is presented in Fig. 5. This difference can be quantified by calculating the $K$ factor, which represents the ratio between the power of the direct ray to the power of the reflected components. The mean value of the $K$ factor along the corridor is 0.32 for the copolar measurements and 0.23 for the crosspolar ones, so the relative contribution of the direct component is more significant in the copolar IR.

The linear dependency between both delay parameters can be measured by their correlation coefficients. These have been calculated and are listed in Table II. The high correlation values suggest that a linear relation in the form $\tau_{\text{corr}}[\text{ns}] = \alpha \tau_{\text{mean}}[\text{ns}] - b$ can model the dependence between
Fig. 6. RMS delay versus mean delay in the corridor: measured data, linear fit, and 90% confidence interval. (a) Copolar and (b) crosspolar component.

Fig. 7. Mean delay and rms delay versus transmitter–receiver separation in the PC laboratory. (a) Copolar and (b) crosspolar component.

Fig. 8. Rms delay versus mean delay in the PC laboratory: measured data, linear fit, and 90% confidence interval. (a) Copolar and (b) crosspolar component.
The results of a linear fit to the measured data are presented in Figs. 6, 8, and 10 and in Table II, including the values of parameters $a$ and $b$, the confidence interval for 90% of the data, and the LMS error of the fit. The goodness of the linear fit has been confirmed by a Kolmogorov–Smirnov test [17], which shows that the residuals follow a Gaussian distribution.

For the three environments considered, it has been found that $\tau_{\text{mean}}$ and $\tau_{\text{rms}}$ are high correlated. It should be noted that both parameters are calculated from the same impulse response and that both may have a dependence on a third common variable that is the Tx–Rx separation. Again, the presence of a dominant component in the copolar IRs makes the linear dependence stronger for this component than for the crosspolar one.
is given. In [16], for an indoor laboratory and in the electronics laboratory, the crosspolar component can be estimated. At each antenna position and the inverse relation between these two parameters. As happens with the rms delay spread, some oscillations are superimposed to the mean coherence bandwidth. It was explained before that these oscillations are caused by the multipath propagation and the absence of a spatial averaging.

The relation between \( B_c \) and \( \tau_{\text{MTS}} \) has been long investigated. A relation of the form \( B_c = C\tau_{\text{MTS}}^{\beta} \), where \( B_c \) is expressed in [MHz] and \( \tau_{\text{MTS}} \) in [ns] [6], has been considered. In order to fit the curve \( B_c = C\tau_{\text{MTS}}^{\beta} \), a log–log transformation of the pairs (\( \tau_{\text{MTS}}, B_c \)) is performed so the relation becomes linear. Then a regression line is fitted to the scatter plot of pairs (\( \ln(\tau_{\text{MTS}}), \ln(B_c) \)). The results of the fit are given in Table III and Figs. 13–15.

The lower bound for \( B_c \) given in [5] has been reported in Figs. 13–15 for comparison purposes. It is observed that most of the pairs (\( \tau_{\text{MTS}}, B_{\text{LOS}} \)) are located above this curve. However, for shorter values of \( \tau_{\text{MTS}} \), there are some pairs below the lower bound. Lower values of \( \tau_{\text{MTS}} \) correspond to shorter separations between antennas. In this situation, the mean signal value changes rapidly and the wide-sense-stationary property assumed in [5] for the lower bound calculation is not verified.

The values for \( C \) and \( \beta \) are lower than those found in [6], where a different frequency band (900–1100 MHz) and correlation level (0.5) have been considered. In the corridor and in the PC laboratory, \( \beta \) around 0.5 for the copolar component and around 0.6 for the crosspolar one. In the electronics laboratory, shorter values of \( \beta \) have been obtained. Due to the dependence of \( \tau_{\text{MTS}} \) and \( B_c \) on the IR shape, the crosspolar component presents a larger decaying slope than the copolar one.

VII. CONCLUSION

The radio channel has been investigated in the 470–862 MHz band in three different indoor environments. A combined method is presented in order to filter noise from the frequency response. With this method, more accurate estimations of \( \tau_{\text{MTS}} \) and \( \tau_{\text{MTS}} \) can be calculated. The effect of four different frequency windows on the results has been analyzed. In general, windows with lower SLLs lead to larger values of \( \tau_{\text{MTS}} \) and \( \tau_{\text{MTS}} \). The Blackman–Harris window yields more accurate estimations and larger values of the delay parameters than other windows.

The variation of \( \tau_{\text{MTS}} \) and \( \tau_{\text{MTS}} \) with the antennas separation has been presented. As the receiver antenna is separated from the transmit antenna, the amplitudes of the reflected signal relative to the direct path become larger, which results in the increase of \( \tau_{\text{MTS}} \) and \( \tau_{\text{MTS}} \). Due to the fact that the copolar component contains the stronger contribution of the LoS wave, the values of \( \tau_{\text{MTS}} \) and \( \tau_{\text{MTS}} \) are lower for the copolar than for the crosspolar component.

It has been demonstrated for the three environments that both parameters exhibit a high correlation and a linear dependence. In all the cases, the correlation and the linear dependence are higher for the copolar than for the crosspolar component. The correlation is also higher in the laboratories than in the corridor. The linear fits have positive slopes that vary with the environment and the polarization component considered.
Fig. 12. $B_c(0.9)$ in the corridor. (a) Copolar and (b) crosspolar component.

Fig. 13. Coherence bandwidth at level 0.9 versus the separation between the transmitter and the receiver in the corridor. (a) Copolar and (b) crosspolar component.

Fig. 14. Coherence bandwidth at level 0.9 versus the separation between the transmitter and the receiver in the PC laboratory. (a) Copolar and (b) crosspolar component.
The coherence bandwidth for a correlation level of 0.9 has been calculated for each of the antenna positions along the corridor. It has been confirmed that an inverse relation between $B_c$ and $\tau_{\text{rms}}$ exists. The results of the fit of the relation with the form $B_c = C\tau_{\text{rms}}^{-\beta}$ are presented. Measurements are also compared with the lower bound given in [5]. The differences in the results for both polarization components are due to the different shape of the IRs.

**REFERENCES**

Mercedes Sánchez Varela was born in Vigo, Spain, in 1974. She received the M.S. degree in telecommunications engineering from the Universidade de Vigo in 1998, where she is currently pursuing the Ph.D. degree in the Departamento de Tecnoloxías das Comunicacións. During 2000, she was a Visiting Researcher at the Department of Engineering Science, University of Oxford, U.K. Her research interests include radio channel characterization and measurement. She has performed measurement campaigns with a vector network analyzer and now is processing results. She is also working on antenna analysis and design for DVB-T distribution networks.

Manuel García Sánchez (S’88–M’93) received the Ingeniero de Telecomunicación degree from the Universidad de Santiago de Compostela, Spain, in 1990 and the Doctor Ingeniero de Telecomunicación degree from the Universidad de Vigo, Spain, in 1996. In 1990, he joined the Departamento de Tecnologías de las Comunicaciones, Universidad de Vigo, where he currently teaches courses in electromagnetic fields as a Profesor Titular de Universidad. He has been a Visiting Researcher at the Departamento de Señales, Sistemas y Radiocomunicaciones, Universidad Politécnica de Madrid, Spain, and the Department of Electronics and Information Technology, University of Glamorgan, U.K. His research interests include studies of indoor and outdoor radio channel modeling for wide-band applications at millimeter-wave frequencies, mobile communications, and DVB-T distribution networks.